Early polarization observations of the optical emission of gamma-ray bursts: GRB 150301B and GRB 150413A

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ABSTRACT
We report early optical linear polarization observations of two gamma-ray bursts made with the MASTER robotic telescope network. We found the minimum polarization for GRB 150301B to be 8 per cent at the beginning of the initial stage, whereas we detected no polarization for GRB 150413A either at the rising branch or after the burst reached the power-law afterglow stage. This is the earliest measurement of the polarization (in cosmological rest frame) of gamma-ray bursts. The primary intent of the paper is to discover optical emission and publish extremely rare (unique) high-quality light curves of the prompt optical emission of gamma-ray bursts during the non-monotonic stage of their evolution. We report that our team has discovered the optical counterpart of one of the bursts, GRB 150413A.

Key words: telescopes – gamma-ray burst: general – gamma-ray burst: individual: GRB 150301B – gamma-ray burst: individual: GRB 150413A.

1 INTRODUCTION
Gamma-ray bursts are among the most distant and powerful explosions in the Universe. The idea that gamma-ray burst emission originates in a narrow magnetized jet with an opening angle of 1° now appears to be a generally accepted hypothesis, though so far unsupported by direct experimental evidence. One of the consequences of the presence of jets and ordered magnetic fields in these jets could be observation of the polarization of the prompt emission of the burst. Prompt optical emission of gamma-ray bursts is the emission that originates at the time when the central engine of the burst still continues to operate, i.e. when the gamma-ray emission has not yet faded.

Scarce and always unique (one burst, one telescope) observations of early optical emission, which is sometimes synchronous with gamma-ray emission, show that bursts may exhibit one of the two behaviour patterns (Vestrand et al. 2005):

(i) optical emission appears simultaneous with gamma-ray emission (or sometimes even before it) and appears correlated with the gamma-ray light curve; and

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of 15–30 deg s\(^{-1}\) and shows no correlation in form with the hard radiation. This dichotomy has been confirmed by later observations of the gamma-ray bursts GRB 100901A and GRB 100906A (Gorbovskoy et al. 2012).

Several authors have reported observing 10–25 per cent polarization of gamma-ray emission (Granot 2003; Yonetoku et al. 2011). In addition, numerous attempts were made in the late 20th and early 21st-century to detect the polarization of the optical afterglow 10–20 h after the burst. It turned out that late afterglow either shows no polarization or only a weak polarization at a level of several percent (Covino et al. 2004). The authors of several early polarization observations of gamma-ray bursts (Mundell et al. 2007, 2013; Steele et al. 2009; Cucchiara et al. 2011; Uehara et al. 2012; King et al. 2014; Kopáč et al. 2015) reported detecting polarization at a 10 per cent level. Also GRB afterglows can be circularly polarized at optical wavelengths (Wiersema et al. 2014). However, so far not a single positive detection has been made of the polarization of the prompt optical emission of the gamma-ray burst. The problem is with the duration of gamma-ray bursts, which are too short and usually last no more than several tens of seconds. Hence, fully automatic observations are needed to be made using robotic telescopes equipped with polarimetry devices. To perform early polarimetric observations of the prompt optical emission of gamma-ray bursts, we developed (Lipunov et al. 2010, 2004) a network of MASTER II robotic telescopes,\(^1\) which we installed in seven sites across the globe (Kornilov et al. 2012). Each telescope has two wide-field (4 deg\(^2\) each) 400-mm diameter optical tubes\(^2\) and is equipped with an extra axis on a superfast mounting capable of moving at a speed of 4 deg\(^{2}\) each second that varies in concert with gamma-ray emission) at the early non monotonic light-curve stage, which then transforms into standard polarization degree of the emission (Pruzhinskaya et al. 2014). The two perpendicular polarization directions yields a lower limit for the polarization plane is inclined by 45° to both of them. The two synchronous frames (taken with different cameras) used to measure \(P_{\text{lowlim}}\) are mutually calibrated, so that the average \(P_{\text{lowlim}}\) for comparison stars would be \(\text{Avg}(P_{\text{lowlim}}) = 0\). This is achieved at the stage of the photometric calibration. We use the same reference stars for frames in both polarizations. The USNO B-1 catalogue that we use is naturally produced in unpolarized light. It means that \(\text{Avg}([\alpha_{\text{left}}] - [\alpha_{\text{right}}]) = 0\) and, consequently \(\text{Avg}(P_{\text{lowlim}}) = 0\). Of course with such a calibration, we zero average polarization of the background stars, i.e., Galactic polarization. However, in the context of this work it is convenient even the object of study is the polarization of extragalactic sources and Galactic polarization must be excluded in any case. Of course this can only work under the assumption of uniform and equal polarization of the background stars. Given the large distance from the galactic plane(galactic latitude (\(b\)) for bursts is \(|b_{\text{GRB150301B}}| = 29.8\) d and \(|b_{\text{GRB150413A}}| = 45.5\) d, not huge Galactic absorption\((A_V = 0.20\) for GRB 150301B and \(A_V = 0.05\) for GRB 150413A (Schlafly & Finkbeiner 2011)) and small area covered by reference stars (0.13 \(\times\) 0.13 deg\(^2\)), we believe it is possible to work in such an approximation.

In the case of the two GRBs considered below, MASTER optical linear polarization observations were made only in two polarizers with no significant difference found in the two channels for any of the events. Unfortunately, it is impossible to make a conclusion about the absence of polarization if observations only in two polarizers are available (the polarization plane may be aligned at 45° relative to the polarizers).

All MASTER magnitudes reported at this paper is unfiltered magnitude calibrated as 0.2\(B + 0.8R\) by the USNO-B1 catalogue. The reasons for using this calibration are discussed in the early papers (Gorbovskoy et al. 2012; Kornilov et al. 2012). MASTER polarization band is determined by the respond curves of the CCD camera and a transmission curve of polarizing filter, which have been reviewed here (Pruzhinskaya et al. 2014).

2 OBSERVATION TECHNIQUE

2.1 Polarimeter

The MASTER net was designed with the objective to deliver polarization information as early as possible after GRB triggers. More than 100 observations of GRBs were made by the MASTER global robotic net. Optical emission was detected for \(\sim 20\) GRBs (Lipunov et al. 2007; Gorbovskoy et al. 2012; Pruzhinskaya et al. 2014). The GRB 100906A, GRB 110422A, and GRB 121011A events deserved attention because their optical observations were carried out during the gamma-ray emission.

We have a linear polarimeters only. So at this article we assume zero circular polarization.

Let \(I_1\) and \(I_2\) be the signals of the detector with orthogonal polarimeter filters (for example 45° and 135° interposed). Then a value \(P_{\text{lowlim}} = (I_1 - I_2)/(I_1 + I_2)\) can be a lower limit of linear polarization.

To derive value of polarization level one needs to perform observations with filters for linear polarization positioned at three angles, for example, at 0°, 45°, and 90° with respect to the reference direction. We have two orthogonal filters at each of the MASTER-II telescopes; thus, if such telescope observes alone, only a lower limit on polarization level – \(P_{\text{lowlim}}\) can be deduced.

If the module of \(P_{\text{lowlim}}\) is less than measurement error \(\sigma_p\), then formally we have zero as a minimum estimate for the degree of linear polarization. 100 per cent linear polarization is also possible in this case, if a filters is accidentally positioned so that a source’s polarization plane is inclined by 45° to both of them.

The two synchronous frames (taken with different cameras) used to measure \(P_{\text{lowlim}}\) are mutually calibrated, so that the average \(P_{\text{lowlim}}\) for comparison stars would be \(\text{Avg}(P_{\text{lowlim}}) = 0\). This is achieved at the stage of the photometric calibration. We use the same reference stars for frames in both polarizations. The USNO B-1 catalogue that we use is naturally produced in unpolarized light. It means that \(\text{Avg}([\alpha_{\text{left}}] - [\alpha_{\text{right}}]) = 0\) and, consequently \(\text{Avg}(P_{\text{lowlim}}) = 0\). Of course with such a calibration, we zero average polarization of the background stars, i.e., Galactic polarization. However, in the context of this work it is convenient even the object of study is the polarization of extragalactic sources and Galactic polarization must be excluded in any case. Of course this can only work under the assumption of uniform and equal polarization of the background stars.

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2.2 Determination of an errors

In the observation of gamma-ray bursts, two factors need to be considered as follows. (a) We needed a high temporal resolution in the initial stages of the light curve. And (b) it is important not to loose a quickly decaying source from view. So MASTER exposure times follow the relation \(t_{\text{exp}} = (T_{\text{start}} - T_{\text{trigger}})/5\), where \(T_{\text{trigger}}\) is
the trigger time (t\textsubscript{\text{trig}}), T\textsubscript{\text{start}} is the time of the beginning of exposure (t). The exposure time is rounded to an integer with a step of 10 s and cannot exceed 3 min. In consequence, frames are quite different from each other and the task of error determination becomes difficult. To do this, we analyse the distribution of P\textsubscript{\text{lowlim}} of the comparison stars, depending on the magnitude. Then, we select only the stars with magnitude in the interval m\textsubscript{\text{grb}} ± 0.5, where m\textsubscript{\text{grb}} is a GRB magnitude on a given frame. Dispersion of comparison stars P\textsubscript{\text{lowlim}} in this interval we call as 1\sigma error of GRB P\textsubscript{\text{lowlim}}. We use the comparison star in a radius of 0.25\arcsec for GRB 150301B and 1\arcsec for GRB 150413A. This difference is due to the bit high value of the absorption in the case of GRB 150301A, and a desire to meet the conditions specified in Section 2.1. Any errors that are reported in this paper is 1\sigma.

### 3 MEASUREMENT OF EARLY POLARIZATION OF THE GAMMA-RAY BURST GRB 150301B

The gamma-ray burst GRB 150301B was recorded and localized by Swift space observatory at 19:38:04 UT on 2015 March 1 (Lien et al. 2015). According to the data from Swift/Burst Alert Telescope (BAT) instrument, GRB 150301B was a long gamma-ray burst with T\textsubscript{\text{je}} = 12.44 ± 1.49 s (Ukwatta et al. 2015).

MASTER-SAAO robotic telescope received the socket message from the GCN network, immediately interrupted its regular survey, pointed to the burst localization area, and started observations in two polarization filters 60 s after notice time and 79 s after trigger time at 2015-03-01 19:39:23 UT. The very first frame showed a bright (15.3 m) optical source with the coordinates

RA(J2000) = 05\textdegree 56\textquoteright 39\arcsec 94

Dec.(J2000) = −57° 58′ 10′′ 0

and an error of 1 arcsecond (Buckley et al. 2015).

The Swift XRT X-ray telescope and UVOT ultraviolet telescope also detected X-ray and optical emission, respectively, from GRB 150301B 82 and 85 s after the BAT trigger (Lien et al. 2015). Later, 5.12 h after the burst, the Very Large Large Telescope (Paranal Observatory, Chile) and X-shooter measured the redshift of this burst, which was found to be z = 1.5169 (de Ugarte Postigo et al. 2015). MASTER observations. MASTER robotic telescope observed this gamma-ray burst in the alert mode in two polarization planes with variable exposures (see Section 2.2).

As a result of our observations, the following light curve was obtained (see Fig. 1 and Table 1).

MASTER-SAAO telescope observations of GRB 150301B lasted from 2015-03-01 19:39:23 UT through 2015-03-01 20:48:11, i.e. for about 1 h beginning with the 79th second after the trigger. Observations were then interrupted because of heavy clouds. We observed well-defined optical emission from GRB 150301B up until 2000 s after the trigger while the source remained brighter than the limiting magnitude on our frames. Unfavorable weather conditions prevented us from extending the light curve by obtaining deep stacked frames. Throughout the entire time interval from 79 to 2000 s the source faded in accordance with the power law F ∝ t^{-1.16 ± 0.05} which, in general, is typical of GRB that time (Mészáros 2006).

#### 3.1 Spectrum from the optical to X-ray

Observations of the gamma-ray burst with MASTER telescopes were performed synchronously with observations on the Swift/XRT X-ray telescope (Fig. 2). Optical data corrected for absorption N\textsubscript{H} = 6.2 × 10\textsuperscript{20} cm\textsuperscript{-2} (Maselli et al. 2015) which is consistent with the Galactic value (Willingale et al. 2013). Note that variability observed in X-rays in the form of three local maxima in the 88, 191 and 282 s does not occur in the optical wavelength, that has already been observed previously (for example, Ziaeepour et al. 2008; Li et al. 2012).

#### 3.2 Polarimetric measurements

The technique of polarimetric measurements used on MASTER network telescopes allows detecting the optical emission polarization at a level of 5 per cent or higher for objects brighter than 16 m or estimating a lower limit for polarization (if the object is observed only by one of the observatories of the network; Pruzhinskaya et al. 2014).

As is evident from the light curve (Fig. 1), the data points for different polarizers differ substantially for the first and fifth observations suggesting possible polarization of the emission.

<table>
<thead>
<tr>
<th>t − T\textsubscript{\text{trig}} (s)</th>
<th>Exp. (s)</th>
<th>Unfiltered P\textsubscript{\text{15}} (m)</th>
<th>Unfiltered P\textsubscript{\text{35}} (m)</th>
</tr>
</thead>
<tbody>
<tr>
<td>79</td>
<td>10</td>
<td>15.26 ± 0.04</td>
<td>15.43 ± 0.05</td>
</tr>
<tr>
<td>127</td>
<td>20</td>
<td>15.61 ± 0.05</td>
<td>15.71 ± 0.05</td>
</tr>
<tr>
<td>187</td>
<td>40</td>
<td>16.17 ± 0.06</td>
<td>16.18 ± 0.06</td>
</tr>
<tr>
<td>276</td>
<td>50</td>
<td>16.47 ± 0.08</td>
<td>16.54 ± 0.07</td>
</tr>
<tr>
<td>389</td>
<td>80</td>
<td>17.14 ± 0.10</td>
<td>16.79 ± 0.08</td>
</tr>
<tr>
<td>502</td>
<td>100</td>
<td>17.48 ± 0.12</td>
<td>17.49 ± 0.10</td>
</tr>
<tr>
<td>659</td>
<td>130</td>
<td>17.89 ± 0.14</td>
<td>17.71 ± 0.10</td>
</tr>
<tr>
<td>827</td>
<td>160</td>
<td>18.13 ± 0.16</td>
<td>18.22 ± 0.12</td>
</tr>
<tr>
<td>1015</td>
<td>180</td>
<td>18.38 ± 0.17</td>
<td>18.16 ± 0.12</td>
</tr>
<tr>
<td>1246</td>
<td>180</td>
<td>18.48 ± 0.18</td>
<td>18.86 ± 0.14</td>
</tr>
<tr>
<td>1465</td>
<td>180</td>
<td>19.01 ± 0.21</td>
<td>18.50 ± 0.13</td>
</tr>
<tr>
<td>1688</td>
<td>180</td>
<td>18.68 ± 0.19</td>
<td>18.85 ± 0.14</td>
</tr>
<tr>
<td>1914</td>
<td>180</td>
<td>18.98 ± 0.21</td>
<td>18.87 ± 0.13</td>
</tr>
<tr>
<td>2142</td>
<td>180</td>
<td>19.03 ± 0.19</td>
<td>19.09 ± 0.14</td>
</tr>
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measurements at different polarizations, which suggests a possible polarization of radiation at a given time. Therefore, for GRB 150301B, we can estimate a lower limit for the polarization degree of its optical emission, which we compute by the following simple formula: $P_{\text{lowlim}} = (I_{\text{90}} - I_{\text{135}})/(I_{\text{90}} + I_{\text{135}})$, where $P_{\text{lowlim}}$ is the lower limit of polarization, i.e. $P_{\text{lowlim}} > 0$, and $I_{\text{90}}$ and $I_{\text{135}}$ are the optical fluxes recorded through a linear polaroid turned by 90° and 135°, respectively. Fig. 3 shows the magnitude dependence of $P_{\text{lowlim}}$ for GRB 150301B and comparison stars. All reference stars were selected within the 256 pixel or 8 arcmin radius from the gamma-ray burst. Stars with cosmetic defects on one of the CCDs (broken and hot pixels) were removed manually. The two synchronous frames (taken with different cameras) used to measure $P_{\text{lowlim}}$ are mutually calibrated, so that the average $P_{\text{lowlim}}$ for comparison stars would be $\text{Avg}(P_{\text{lowlim}}) = 0$. To determine the error of the inferred lower polarization limit for GRB 150301, we measured the standard deviation of $P_{\text{lowlim}}$ from zero for the comparison stars in the magnitude interval $M_{\text{GRB150301B}} \pm 0.5$, i.e. from 14.9 to 15.9.

The lower limit for the polarization of GRB 150301 during the time interval from 79 to 89 s is thus equal to $P_{\text{GRB150301B}} > 7.57 \pm 0.04$ at a 3.2σ significance level. Despite the fact that formally the polarization of the 5th point (400 s after the burst) is great, but the object is visible almost at the limit of detection as shown in Fig. 4 can not find a suitable comparison star magnitudes.

4 OBSERVATIONS OF GRB 150413A

This gamma-ray burst was detected by BAT instrument of Swift observatory (Markwardt et al. 2015). At 13:54:58 UT, the Swift BAT triggered and located GRB 150413A. Unfortunately, due to an observing constraint X-ray observations of GRB 150413A were not made.

The telescope of MASTER robotic network located at Tunka astrophysical centre near Baikal lake was automatically pointed to the target 50 s after receiving the alert message and 132 s after the trigger operated onboard SWIFT space observatory.

We found nothing new in the first frame within the error box of SWIFT BAT telescope. However, on the second and next frames a hitherto unknown source was found by MASTER optical transient autodetection system at a position $\text{RA}(J2000) = 12^h41^m41^s.98$ Dec.$(J2000) = +71^\circ50^\prime28^\prime$
and with uncertainty 1 arcsec and unfiltered magnitude about $m = 16.1$ m (Ivanov et al. 2015). (Fig. 5)

Its brightness increased steadily and by the 400th second the object became brighter than 15 m Tyurina et al. (2015). We seemed to have found an unusual gamma-ray burst, which remained brighter for longer than in most previous events.

However, following the MASTER detection, the 2.2 m telescope at the German–Spanish Calar Alto (CAHA) Observatory observed the afterglow starting at 20:25 UT. (i.e. 6.5 h postburst) under poor weather and seeing conditions (Sanchez-Ramirez et al. 2015). Data was reduced and calibrated on the usual way using IRAF, and custom tools coded up in python. The resulting spectrum has a low SNR ratio, but we clearly detect a broad absorption line which we identify as Lyman-alpha. No metal line is significantly detected in the wavelength range covered due to the observing conditions. We performed a Voigt profile fitting using Ly $\alpha$ and Ly $\beta$ lines, finding a value of $z = 3.140 \pm 0.005$ for the redshift and $\log(N_H) = 22.0 \pm 0.25$ for the HI column density (Fig. 6). This redshift is well constrained due to the identification of the Ly $\beta$ line, and consistent with the later provided by de Ugarte Postigo and Tomasella (de Ugarte Postigo & Tomasella 2015). The value of the HI column density we find is hardly seen in QSO sight lines, but quite common in GRB-DLAs (e.g. Cucchiara et al. 2015).

So everything became clear. If the gamma-ray burst had a redshift $\sim 3.1$ (Sanchez-Ramirez et al. 2015; de Ugarte Postigo & Tomasella 2015), it means that all processes appear slowed down by a factor of $(1 + z)^4$. With this correction applied, we converted our observations to the comoving reference frame of the object and found the event to have a proper duration of less than 1 min, and hence to be a rather typical long gamma-ray burst most likely resulting from the formation of a black hole after the core collapse of a massive star.

The light curve of GRB 150413A shown in Fig. 7. This gamma-ray burst, as the previous one was observed only by one telescope of MASTER net–MASTER-Tunka. The observations at the telescope MASTER-Tunka were made in two polarization angles (0 and 90°, respectively). Therefore, as in the previous case, we can talk about the lower limit of the possible polarization of the afterglow of GRB 150413A.

Fig. 8 is a graph of the degree of polarization of GRB and background stars in the first frame. Diagram plot on the assumption that all the brightest stars (brighter than 15 mag) in the frame have zero polarization. Following this assumption, the median value of $P_{\text{Polar}} = (I_1 - I_2)/(I_1 + I_2)$ sample of stars brighter than 15 mag were normalized to zero. In general (for more than 95 per cent of the stars), the diagram keeps symmetry up to the limit at about 17 m, with increasing spread is inversely proportional to the magnitude of star. This spread will characterize the measurement error. Huge scatter
Figure 8. Value of the polarization low limit ($P_{\text{lowlim}}$) as a function of the magnitude on the second frame of GRB 150413A observations. Explanations asymmetry present image and fig. 8 given in chapter 4.

Figure 9. Value of the polarization low limit ($P_{\text{lowlim}}$) as a function of the magnitude on the first five frames. All GRB points covered by reference stars. So there is no evidence to believe that there is a polarization for this GRB for the bright stars (on the level of polarization like 40 per cent) related to the false identification of closely spaced stars. Strong variations in $m > 17$ m region 0.25 connected with small difference of the upper limits of two frames in different polarizations.

As can be seen from Fig. 8, despite formally a significant degree of polarization of the first measured point ($P_{\text{lowlim}}$ 15 per cent), it is less than the measurement error. That is not surprising since the burst in this picture has a very close to limiting magnitude. On all subsequent (till a maximum burst luminosity) frames polarization also not found, because all polarization measurements fall within the basic distribution of background stars, as seen on Fig. 9. It shows all the background stars with an average polarization for the first five images and points represent gamma-ray burst on these frames. Diagram is plotted on the same principle as the previous one.

5 DISCUSSION

As we already pointed out, (Vestrand et al. 2005) noted that two behaviour types of early optical emission are observed for gamma-ray bursts. In some cases, optical emission varies in concert with gamma-ray flux, whereas in other cases the optical light curve is totally uncorrelated with the gamma-ray flux and usually consists of a smooth rise followed by power-law afterglow. We earlier observed both types of such behaviour (Gorbovskoy et al. 2012). It is logical to suggest (Vestrand et al. 2005) that in the former case we are dealing with the optical emission arising in the internal shock region, where it is generated simultaneously with the gamma-ray emission. In this case, the optical light curve reflects the internal properties of the relativistic magnetized jet (ultrarelativistic ejecta driven by the GRB) and detection of its polarization would provide unique information about the jet structure. The optical emission of the second type may be associated with a different region of the jet (in the standard fireball model, Mészáros & Rees 1999), which forms as a result of relativistic jet colliding with interstellar gas compressed in the bow shock or with the stellar wind of the progenitor.

The gamma-ray bursts that we report in this paper belong to the second type. It is for such bursts that earlier polarization was observed at the afterglow stage. Our observations very much confuse the problem of the optical emission polarization of gamma-ray bursts. The point is that we detected polarization only for the GRB 150301B burst and only at certain rather early time instants. On a diagram (Fig. 10), we plot MASTER GRB 150301B polarization low limit in comparison with other known earlier measurements of the GRB polarization. At the moment, it is the earliest measurement (from among successful) optical linear polarization. The observation is in agreement with other measurements and shows the ability to burst to have a high degree of optical polarization in the earliest time.

We plan to defer the theoretical and phenomenological description of the nature of prompt optical emission of gamma-ray bursts and its polarization by several years, after sufficient amount of positive polarimetric observations of gamma-ray bursts becomes available. Our optimism is inspired by the doubling of the number of MASTER network telescopes deployed in 2015 in such sites of excellent astroclimate as Canary islands, South Africa, and Argentina.

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